MEASUREMENT OF THE POLYTROPY INDEX FOR GAS-DETONATION PRODUCTS

## Yu. N. Denisov and Ya. K. Troshín

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This method for  $\gamma$  is based on measurement of  $D_2/D_1$  (ratio of velocities of incident and reflected waves) for two colliding detonation waves under conditions where one can neglect the three-dimensional structure of the reaction zone in the incident wave. The collision of these effectively one-dimensional detonation waves is described by the solution for the collision of a detonation wave with an absolutely rigid wall for the strong-wave approximation [1,2]. This approximation may give a substantial error in  $\gamma$  for a gas detonation, where the ratio of  $p_0$  (initial pressure) to  $p_1$  (pressure in detonation wave) is 0.05-0.2. We therefore deduce the relation of  $\gamma$  to  $D_2/D_1$  without neglecting  $p_0$  relative to  $p_1$ .

The conservation and other equations for the incident wave are

$$\rho_0 D_1 = \rho_1 (D_1 - u_1), \qquad p_1 - p_0 = \rho_0 D_1 u_1,$$
$$\frac{\rho_1}{\rho_0} = \frac{\gamma_1 + 1 - p_0 / p_1}{\gamma_1}. \tag{1}$$

For the reflected wave (on the assumption that  $\gamma_1 = \gamma_2 = \gamma$ ) the conservation laws give

$$\rho_2 D_2 = \rho_1 (D_2 + u_1), \quad p_2 - p_1 = \rho_1 (D_2 + u_1) u_1,$$
$$\frac{\rho_2}{\rho_1} = \frac{p_1 (\gamma - 1) + p_2 (\gamma + 1)}{p_1 (\gamma + 1) + p_2 (\gamma - 1)}.$$
(2)

Here  $\rho$  and u are respectively the density and mass velocity, while subscripts 0, 1, and 2 relate to the initial mixture, the incident wave,



Fig. 1. Relation of  $\gamma$  to  $D_2/D_1$ .

and the reflected wave. In Eqs. (1) and (2) we make the substitutions

$$d = \frac{D_2}{D_1} + 1, \quad \xi = \frac{\gamma + 1}{\gamma}, \quad \pi = \frac{p_0}{p_1}$$
 (3)

to get

$$l\left(1+\frac{1-\pi}{\gamma}\right)-1-\frac{\xi(1-\pi)}{4}-\left(\frac{\xi^{2}(1-\pi)^{2}}{16}+1\right)^{1/2}=0.$$
 (4)

This may be put as

$$a\gamma^2 + b\gamma + c = 0. \tag{5}$$

If the terms in  $\pi^2$  are neglected, the coefficients are

$$a = 2d^{2} - 5d + 1 + \pi (d - 1),$$
  

$$b = 4d^{2} - 6d + 1 + \pi (7d - 4d^{2} - 1),$$
  

$$c = d (2d - 1) - \pi (2d - 1) d.$$
(6)

Figure 1 shows solutions to Eq. (5) for  $\gamma$  positive and for the most probable range of  $D_2/D_1$  (0.3-0.5) for various  $\pi = p_0/p_1$ : the upper curve corresponds to the solution for a strong detonation wave.



Fig. 2. The detonation tube: 1) metal tube for initiating detonation, 2) wire spiral for igniting gas mixture, 3) Shchekin spiral, 4) metal tube, 5) experimental section,

This may be used as a nomogram with the observed  $D_2/D_1$  and  $p_0/p_1$  to deduce  $\gamma$ .

Figure 2 shows the apparatus used to measure  $\gamma$ . The spiral 2 ignites the gas mixture in tube 1; the conversion of the combustion to a detonation is accelerated [3] by the spiral 3. Tube 1 is connected to the symmetrically placed tube 4; part 5 is made of glass and has an internal diameter of 16 mm. The middle part is viewed by a ZhFR-1 camera [4, 5].

The gas mixtures had compositions and initial pressures far from the detonation limits. The recordings were as shown in Fig. 3;  $D_2/D_1$ is deduced from the angles  $\alpha$  and  $\beta$  as follows:

$$\frac{D_2}{D_1} \coloneqq \frac{\operatorname{tg} \frac{1}{2\beta}}{\operatorname{tg} \frac{1}{2\alpha}} \,.$$

Table 1 gives results for various mixtures, including acetyleneoxygen containing argon, for which  $\gamma = 1.67$ , it being assumed that the  $\gamma$  for this case would be larger than that for a mixture without argon. The  $p_0/p_1$  for these mixtures were as follows:  $2H_2 + O_2 0.05$ ,  $CH_4 + 2O_2 0.03$ ,  $C_2H_2 + 2.5O_2 0.03$ . For  $p_0$  of 500 mm Hg or less, the mixtures  $2H_2 + O_2$ ,  $CH_4 + 2O_2$ , and  $C_2H_2 + 2.5O_2 + 2.5Ar$  showed an increase in  $D_2/D_1$  as  $p_0$  decreased, which evidently reflects the influence of the wave structure [6-8]. Inhomogeneity produces a complex pattern in the reflected waves near the point of reflection, the lines of propagation of the fronts being bent (Fig. 3a). There is hardly any effect on  $D_2/D_1$  for  $p_0 > 500$  mm Hg, and so  $\gamma$  was deduced from the

Table 1						
$D_2/D_1$ for Various	Initial	Pressures	p <sub>0</sub> in	Gas	Mixtures	

Pa	$D_{s} D_{1}$					
	$2H_2 + O_2$	CH4 + 2O2	$C_{2}H_{2} + 2.5O_{2}$	$C_2H_2 + 2.5O_2 + 2.5Ar$		
915 880 760 500	$\begin{array}{c} 0.42 \pm 0.04 \\ 0.42 \pm 0.01 \\ 0.45 \pm 0.02 \end{array}$	$0.422 \pm 0.005$ 0.420 \pm 0.02 0.424 \pm 0.02	$\begin{array}{ c c c c c c c c c c c c c c c c c c c$	$\begin{array}{ c c c c c c c c c c c c c c c c c c c$		





Fig. 3. Recordings of colliding detonation waves: a)  $2H_2 + O_2$ ,  $p_0 = 880$  mm Hg, 16 mm tube; b)  $C_2H_2 + 2.5O_2$ ,  $p_0 = 300$  mm Hg, 16 mm tube.

Mixture	$2H_2 + O_2$	$CH_4 + 2O_2$	$C_2H_2 + 2.5O_2$	$C_2H_2 + 2.5O_2 + 2.5Ar$
r	$1.225 \pm 0.025$	$1.245 \pm 0.015$	$1.16 \pm 0.03$	$1.225 \pm 0.022$

data of Table 1 only for  $\rho_0 > 500$  mm Hg, for which Fig. 1 gives the  $\gamma$  listed in Table 2.

It is clear that argon does increase  $\gamma$  for acetylene-oxygen mixtures.

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